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## Magnetically frustrated synthetic end member $\text{Mn}_2(\text{PO}_4)\text{OH}$ in the triplite–triploidite family†

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The manganese end member of triplite–triploidite series of compounds,  $\text{Mn}_2(\text{PO}_4)\text{OH}$ , is synthesized by a hydrothermal method. Its crystal structure is refined in the space group  $P2_1/c$  with  $a = 12.411(1)$  Å,  $b = 13.323(1)$  Å,  $c = 10.014(1)$  Å,  $\beta = 108.16(1)^\circ$ ,  $V = 1573.3$  Å<sup>3</sup>,  $Z = 8$ , and  $R = 0.0375$ . Evidenced in measurements of magnetization  $M$  and specific heat  $C_p$ ,  $\text{Mn}_2(\text{PO}_4)\text{OH}$  reaches a long range antiferromagnetic order at  $T_N = 4.6$  K. As opposed to both triplite  $\text{Mn}_2(\text{PO}_4)\text{F}$  and triploidite-type  $\text{Co}_2(\text{PO}_4)\text{F}$ , the title compound is magnetically frustrated being characterized by the ratio of Curie–Weiss temperature  $\theta$  to Néel temperature  $T_N$  of about 20. The large value of frustration strength  $|\theta|/T_N$  stems from the twisted saw tooth chain geometry of corner sharing triangles of Mn polyhedra, which may be isolated within tubular fragments of a triploidite crystal structure.

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## Introduction

The key factors leading to the suppression of the magnetic ordering temperature as compared to the values of leading exchange interactions in the crystal are reduced dimensionality and frustration. The combination of these factors may result in huge values of the ratio  $|\theta|/T_N \geq 50$  realized in some triangular or kagomé lattices,<sup>1,2</sup> albeit some recently investigated representatives of this multitude possess a frustration strength  $|\theta|/T_N \leq 10$ .<sup>3</sup> Comparable values of the frustration strength are realized in some three-dimensional magnets represented by numerous garnets, pyrochlores or spinels.<sup>4–6</sup> Nevertheless, any object with Curie–Weiss temperature  $\theta$  exceeding Néel temperature  $T_N$  by an order of magnitude belongs to the field of highly frustrated magnetism. This flourishing field attracts attention as an “excellent playground to discover new states and new properties of matter” including gapped and gapless spin liquids, spin ices, spin nematics, *etc.*<sup>7</sup> It should be noted, however, that these ground states should be distinguished from those realized in magnetically ordered, commensurate or

incommensurate, magnetic structures due to some degree of frustration.

The geometric frustration in three-dimensional magnets is associated usually with the triangular motif in the arrangement of magnetically active ions. Thus, a three-dimensional network of corner-sharing tetrahedra constitutes the magnetic subsystem in cubic transition metal spinels and in rare-earth pyrochlores. The ferrite structure can be considered as a sequence of frustrated kagomé and triangular layers stacked alternatively to form a three-dimensional arrangement. Significantly more complicated is the magnetic network of frustrating units in garnets represented by interpenetrating rings of corner-sharing triangles. The transition metal based members of the garnet’s family exhibit only modest degree of frustration  $|\theta|/T_N \sim 3$ .<sup>8–10</sup> Also, there are three-dimensional heavily frustrated magnets having hyper-kagomé<sup>11</sup> or tetragonal pyrogermanate structures,<sup>12</sup> but not many.

In this work, we present a candidate for a frustrated magnet that belongs to a completely different structure type, *i.e.* triploidite  $\text{Mn}_2(\text{PO}_4)\text{OH}$ . The name “triploidite” indicates “near triplite” and specifies a manganese phosphate similar to triplite,  $\text{Mn}_2(\text{PO}_4)\text{F}$ , but containing hydroxyl instead of fluorine. A triploidite group of minerals unites isotypic monoclinic phosphates and arsenates with the general formula  $\text{M}_2(\text{TO}_4)\text{Z}$  and besides triploidite itself,  $\text{Mn}_2(\text{PO}_4)(\text{OH})$ , includes wagnerite,  $\text{Mg}_2(\text{PO}_4)\text{F}$ , hydroxylwagnerite,  $\text{Mg}_2(\text{PO}_4)(\text{OH})$ , wolfeite,  $\text{Fe}_2(\text{PO}_4)(\text{OH})$  and sarkinite,  $\text{Mn}_2(\text{AsO}_4)(\text{OH})$ . All of them crystallize in a space group  $P2_1/c$  and have a  $b$  cell parameter of about 13 Å. For triplite group of minerals described by the same general formula, the  $b$  cell parameter is twice smaller and the monoclinic unit cell is Bravais centered (space group

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$C2/c$ ).<sup>13</sup> Nevertheless, the crystal structures of minerals of both groups have very similar topologies. Natural mineral phases belonging to the triplite–triploidite supergroup<sup>14</sup> often present solid solutions of Mg/Fe/Mn on one hand and F/OH on the other hand. Recent X-ray diffraction studies of more than 40 mineral samples have shown that their crystal structures may be periodically or aperiodically modulated. It has been suggested that they should be described in the framework of a superspace model with the  $(3 + 1)$ -dimensional space group  $C2/c(0\beta 0)s0$  based on the average triplite structure and the modulation vector  $\mathbf{q} = \beta \mathbf{b}^*$ .<sup>15</sup>

Besides minerals, the large series of synthetic analogues – mainly phosphates and some arsenates – were obtained and structurally investigated (Table S1†). They are the end members without isomorphic substitutions and thus, no modulations have been noticed in connection with these triplite- or triploidite-like crystal structures. Curiously, synthetic phases with triplite structures present fluoride-phosphate solely, while triploidite-like compounds are possible to grow both as F- or OH-varieties. Despite the fact that many chemical compositions were realized as laboratory prototypes of minerals, the synthetic triploidite  $\text{Mn}_2(\text{PO}_4)\text{OH}$  has not been obtained so far. Here, we report the hydrothermal synthesis, crystal structure refinement and thermodynamic properties of  $\text{Mn}_2(\text{PO}_4)\text{OH}$ . Unexpectedly, this compound was found to be magnetically frustrated, as opposed to both triplite  $\text{Mn}_2(\text{PO}_4)\text{F}$  and triploidite-type  $\text{Co}_2(\text{PO}_4)\text{F}$ , but similar to  $\text{Mn}_2(\text{AsO}_4)\text{OH}$ .<sup>16</sup> In variance with numerous frustrated antiferromagnets of a pyrochlore or spinel type possessing a three-dimensional triangular magnetic lattice, the  $\text{Mn}_2(\text{PO}_4)\text{OH}$  represents a rare case of corner-sharing twisted saw tooth chain magnetic system.

## Experimental

Single crystals of  $\text{Mn}_2(\text{PO}_4)\text{OH}$  were synthesized under mild hydrothermal conditions. A starting mixture of  $\text{MnCl}_2\text{--H}_3\text{PO}_4\text{--Na}_2\text{B}_4\text{O}_7$  in a 1:1:1 weight ratio was dissolved in distilled water, sealed in a PTFE-lined stainless-steel pressure vessel (fill factor 80%) and kept at 170 °C for 18 days, followed by slow cooling to room temperature. The reaction products were washed with water and dried. The colourless transparent needle-like single crystals of  $\text{Mn}_2(\text{PO}_4)(\text{OH})$  of about 600  $\mu\text{m}$  in length are shown in Fig. 1. These crystals constitute about 70% of the total yield along with fine-crystalline ameghinite,  $\text{NaB}_3\text{O}_5(\text{H}_2\text{O})_2$ , as confirmed by X-ray diffraction analysis of a powder sample.

The composition of the title compound was analyzed with a scanning electron microscope JSM6480LV equipped with an INCA Energy-350 energy dispersive detector and an INCAWave-500 four-crystal wavelength dispersive spectrometer.‡ The measurements were made at 20 kV and 7 nA, and the sample



Fig. 1 Photograph of crystals of the title compound.

was stable under these conditions. X-ray spectral analysis provided a semi-quantitative result with the Mn : P : O ratios close to 2 : 1 : 5, which is consistent with the results of the X-ray diffraction structural study. No impurity elements were found in the crystals' composition.

Single crystal X-ray diffraction data were collected at ambient temperature by using  $\text{MoK}_\alpha$  radiation with an Oxford Diffraction Xcalibur-S area detector diffractometer. Recorded reflection intensities (full sphere of reciprocal space) were corrected for Lorentz and polarization effects, and a numerical absorption correction based on Gaussian integration over a multifaceted crystal model was applied. All calculations were performed in the WinGX32 software package.<sup>17</sup> Atomic scattering factors and anomalous dispersion corrections were taken from the International Tables for Crystallography.<sup>18</sup> The crystal structure was solved *via* direct methods in the space group  $P2_1/c$  with the SHELXT program.<sup>19</sup> Atomic positions were refined against the  $F^2$  data using SHELXL programs<sup>20</sup> with anisotropic displacement parameters for all non-H atoms. Four symmetrically independent H atoms forming hydroxyl groups were obtained by difference-Fourier techniques and refined in an isotropic approximation. The O–H bond lengths were fixed by hard restraints to an empirical value of 0.80 Å in order to obtain comparable H-bond geometry not affected by arbitrary scattering of the refined O–H distances. The main crystallographic characteristics of synthetic triploidite  $\text{Mn}_2(\text{PO}_4)\text{OH}$  and experimental parameters for data collection and structure refinement are summarized in Table 1.

The final atom positions are given in Table S2,† while Table 2 reports characteristic distances. The hydrogen bonds' geometry is presented in Table 3. A bond-valence calculation (Table S3†) has been performed using the algorithm and parameters given in ref. 21. Valence contributions of the H atoms were estimated from O–O distances following the equations from ref. 22; the relationship between the O–O distance and H–O (acceptor) bond valence is calculated from the O–H curve of ref. 23. Data from Table S3† confirm clearly the assignment of O and OH ligands; these data are consistent with the assumed oxidation state of manganese ions, *i.e.*  $\text{Mn}^{2+}$ .

‡ The analysis was performed at the Laboratory of Local Methods for Matter Investigation of the Department of Petrology at the Faculty of Geology of Moscow State University.



**Table 1** Crystal information and details of data collection and refinement

Formula	Mn <sub>2</sub> (PO <sub>4</sub> )(OH)
Absorption $\mu$ [mm <sup>-1</sup> ]	6.725
Space group	<i>P</i> 2 <sub>1</sub> / <i>c</i> , <i>Z</i> = 8
<i>a</i> , <i>b</i> , <i>c</i> (Å)	10.0142(8), 13.323(1), 12.411(1)
$\beta$ (°), <i>V</i> (Å <sup>3</sup> )	108.16(1), 1573.3(1)
<i>F</i> (000), density <i>D</i> <sub>calc</sub> [g cm <sup>-3</sup> ]	1696, 3.747
Crystal size [mm]	0.08 × 0.01 × 0.01
Radiation	MoK $\alpha$ ( $\lambda$ = 0.71073 Å)
Temperature [K]	293(2)
Scanning mode	$\omega$ scan
Measuring range	Max $\theta$ = 26.50°
<i>h</i> , <i>k</i> , <i>l</i> ranges	−12 ≤ <i>h</i> ≤ 12, −16 ≤ <i>k</i> ≤ 16, −15 ≤ <i>l</i> ≤ 15
Absorption correction	Numerical (Gaussian)
<i>T</i> <sub>min</sub> , <i>T</i> <sub>max</sub>	0.857, 0.966
<i>R</i> <sub>int</sub> / <i>R</i> <sub><math>\sigma</math></sub>	0.043/0.023
Reflections total/unique/observed	23 812/3261/2701
[ <i>I</i> > 2 $\sigma$ ( <i>I</i> )]	
Parameters	306
<i>R</i> (observed refl.)	0.0375
<i>R</i> , <i>wR</i> <sub>2</sub> (all refl.)	0.0510, 0.0668
Goodness of fit	1.295
Extinction coefficient	0.00035(5)
$\delta\rho$ (max)/(min) [e Å <sup>-3</sup> ]	0.491/−0.462

**Table 2** Mn<sub>2</sub>(PO<sub>4</sub>)OH bond lengths, Å

Mn1-octahedron	Mn2-bipyramid	Mn3-bipyramid
Mn1–O12 2.139(3)	Mn2–O12 <sup>#1</sup> 2.087(3)	Mn3–O9 2.083(3)
–O4 2.158(3)	–O19 2.099(2)	–O17 2.101(3)
–O6 2.170(3)	–O14 2.116(3)	–O15 2.102(3)
–O16 2.202(2)	–O7 2.134(3)	–O18 2.109(3)
–O5 2.209(3)	–O4 2.189(3)	–O17 <sup>#2</sup> 2.223(3)
–O7 2.299(3)		
<Mn1–O> 2.196	<Mn2–O> 2.125	<Mn3–O> 2.124
Mn4-octahedron	Mn5-octahedron	Mn6-bipyramid
Mn4–O9 2.150(3)	Mn5–O12 <sup>#4</sup> 2.166(3)	Mn6–O5 2.086(3)
–O3 2.185(3)	–O13 2.166(3)	–O10 2.109(3)
–O2 2.194(3)	–O2 2.184(3)	–O18 2.126(3)
–O8 2.196(3)	–O16 <sup>#4</sup> 2.234(3)	–O19 <sup>#5</sup> 2.204(3)
–O1 <sup>#3</sup> 2.237(3)	–O8 2.258(3)	–O20 2.239(3)
–O1 2.239(3)	–O11 2.340(3)	
<Mn4–O> 2.200	<Mn5–O> 2.225	<Mn6–O> 2.153
Mn7-bipyramid	Mn8-octahedron	P1-tetrahedron
Mn7–O8 2.095(3)	Mn8–O9 2.174(3)	P1–O11 <sup>#8</sup> 1.534(3)
–O20 <sup>#6</sup> 2.130(3)	–O6 2.180(3)	–O15 <sup>#9</sup> 1.540(3)
–O14 <sup>#7</sup> 2.154(3)	–O11 <sup>#6</sup> 2.194(3)	–O14 <sup>#10</sup> 1.547(3)
–O15 2.181(3)	–O5 2.229(3)	–O7 1.548(3)
–O10 <sup>#6</sup> 2.188(3)	–O13 <sup>#6</sup> 2.275(3)	<P1–O> 1.542
	–O3 2.320(3)	
<Mn7–O> 2.150	<Mn8–O> 2.229	
P2-tetrahedron	P3-tetrahedron	P4-tetrahedron
P2–O16 <sup>#4</sup> 1.535(3)	P3–O17 <sup>#1</sup> 1.531(3)	P4–O3 1.533(3)
–O13 <sup>#1</sup> 1.542(3)	–O2 1.536(3)	–O6 <sup>#11</sup> 1.537(3)
–O18 <sup>#1</sup> 1.550(3)	–O19 1.542(3)	–O4 1.545(3)
–O1 1.553(3)	–O20 1.546(3)	–O10 <sup>#1</sup> 1.554(3)
<P2–O> 1.545	<P3–O> 1.539	<P4–O> 1.542

Symmetry transformations used to generate equivalent atoms: #1 *x*, −*y* +  $\frac{1}{2}$ , *z* +  $\frac{1}{2}$ ; #2 −*x*, −*y*, −*z* + 1; #3 −*x*, −*y*, −*z* + 2; #4 *x* − 1, *y*, *z*; #5 *x*, −*y* +  $\frac{1}{2}$ , *z* −  $\frac{1}{2}$ ; #6 −*x*, *y* −  $\frac{1}{2}$ , −*z* +  $\frac{1}{2}$ ; #7 *x* − 1, −*y* +  $\frac{1}{2}$ , *z* −  $\frac{1}{2}$ ; #8 *x* + 1, *y*, *z*; #9 *x* + 1, −*y* +  $\frac{1}{2}$ , *z* +  $\frac{1}{2}$ ; #10 −*x* + 1, −*y* + 1, −*z* + 2; #11 −*x* + 1, −*y*, −*z* + 2; #12 −*x*, *y* +  $\frac{1}{2}$ , −*z* +  $\frac{1}{2}$ .

**Table 3** Geometry of hydrogen bonds

D–H...A	<i>d</i> (D–H), Å	<i>d</i> (H...A), Å	<i>d</i> (D...A), Å	$\angle$ (DHA), °
O9–H1...O13 <sup>#6</sup>	0.800(1)	2.24(4)	2.784(4)	125(4)
O9–H1...O17 <sup>#2</sup>	0.800(1)	2.33(4)	2.811(4)	120(4)
O5–H2...O3	0.800(1)	2.36(3)	2.872(4)	122(4)
O5–H2...O4	0.800(1)	2.56(3)	3.003(4)	117(3)
O12–H3...O4 <sup>#5</sup>	0.800(1)	2.36(4)	2.818(4)	117(4)
O12–H3...O10	0.800(1)	2.60(3)	3.182(4)	131(4)
O12–H3...O11 <sup>#8</sup>	0.800(1)	2.34(4)	2.885(4)	126(4)
O8–H4...O1	0.800(1)	2.65(3)	3.122(4)	120(3)
O8–H4...O16 <sup>#4</sup>	0.800(1)	2.23(4)	2.767(4)	125(4)

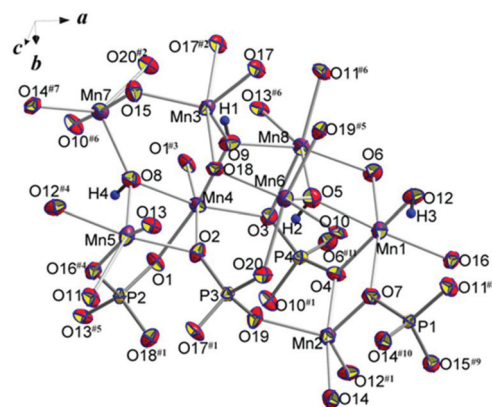
Symmetry transformations used to generate equivalent atoms are as shown in footnotes for Table 2.

## Results

### Crystal structure

The basic structural units of synthetic triploidite, Mn<sub>2</sub>(PO<sub>4</sub>)OH, are shown in Fig. 2. Eight symmetrically independent Mn<sup>2+</sup> ions, all in general structural positions are surrounded by six or five oxygen atoms forming octahedra and five-vertex polyhedra. The interatomic Mn–O distances in the Mn-centered octahedra vary from 2.139(3) to 2.340(3) Å. Each of the four MnO<sub>4</sub>(OH)<sub>2</sub> octahedra has in its first coordination sphere two OH ligands in the *cif*-conformation. The morphology of MnO<sub>4</sub>(OH) five-vertex polyhedra can be defined as distorted trigonal-bipyramids with Mn–O distances ranging between 2.083(3) and 2.239(3) Å. The shortest bond lengths match Mn–OH contacts in all bipyramids (Table 2). The four non-equivalent P atoms have similarly distorted tetrahedral surroundings with P–O distances that vary from 1.531(3) to 1.554(3) Å, and nearly equal average P–O values of 1.54 Å.

The positions and refinement of the H atoms has allowed a detailed analysis of the H bonds' peculiarities (Table 3). The system of O–H–O asymmetrical H bonds between OH groups and O atoms provides additional crosslinking in the structure. With O–O distances between 2.767(4) and 3.182(4) Å, they are classified as weak and very weak bonds. Three of the four



**Fig. 2** Basic structural units with atom labeling scheme. Displacement ellipsoids are presented at the 90% probability level. See symmetry codes in footnotes for Table 2.





hydroxyl groups are involved in bifurcated hydrogen bonding, while the fourth hydrogen atom (H3) forms a trifurcated hydrogen bond. Typically for inorganic compounds, all hydrogen bonds are meaningfully non-linear with O–H...O angles ranging between 117(3) and 126(4)°.

Three varieties of topologically distinct strongly corrugated chains built of Mn-centered polyhedra sharing edges and vertices may be distinguished in the crystal structure, as shown in Fig. 3. The chains of the first type are aligned parallel to the *b* axis of the unit cell and are formed from alternating pairs of sharing edges of MnO<sub>4</sub>(OH)<sub>2</sub> octahedra and MnO<sub>4</sub>(OH) bipyramids linked through common vertices of neighbouring polyhedra according to the scheme: Mn5O<sub>4</sub>(OH)<sub>2</sub> = Mn8O<sub>4</sub>(OH)<sub>2</sub>–Mn6O<sub>4</sub>(OH) = Mn7O<sub>4</sub>(OH)–Mn5O<sub>4</sub>(OH)<sub>2</sub> = Mn8O<sub>4</sub>(OH)<sub>2</sub>..., where “=” denotes the shared edges and “–” signifies the shared polyhedron vertices. The same way of alternation (octahedron = octahedron–bipyramid = bipyramid–octahedron = octahedron...) is realized in the chains of the second type ranged along the *c* axis and built of Mn-centered polyhedra sharing edges and vertices: Mn4O<sub>4</sub>(OH)<sub>2</sub> = Mn4O<sub>4</sub>(OH)<sub>2</sub>–Mn3O<sub>4</sub>(OH) = Mn3O<sub>4</sub>(OH)–Mn4O<sub>4</sub>(OH)<sub>2</sub> = Mn4O<sub>4</sub>(OH)<sub>2</sub>.... The third sort of chains also parallel to the [001] direction, consists of Mn2- and Mn1-centered polyhedra, which are placed along the chain in the following mode: Mn2O<sub>4</sub>(OH) = Mn1O<sub>4</sub>(OH)<sub>2</sub>–Mn2O<sub>4</sub>(OH) = Mn1O<sub>4</sub>(OH)<sub>2</sub>–Mn2O<sub>4</sub>(OH) = Mn1O<sub>4</sub>(OH)<sub>2</sub>..., e.g. the octahedra and trigonal bipyramids share edges and vertices, and alternate within this chain in contrast to the case of triplite where an interchange of polyhedra obeys the rule: two octahedra–two bipyramids. Along the *a* axis the described chains of octahedra and trigonal bipyramids interconnect in a three-dimensional cationic framework through common edges and vertices as shown in Fig. 4. Anionic [PO<sub>4</sub>]<sup>3–</sup> oxo-complexes strengthen the framework by sharing all vertices of tetrahedra with Mn-centered polyhedra (Fig. S1†).

From the magnetic standpoint one can presume that edge-sharing polyhedra are coupled by ferromagnetic exchange

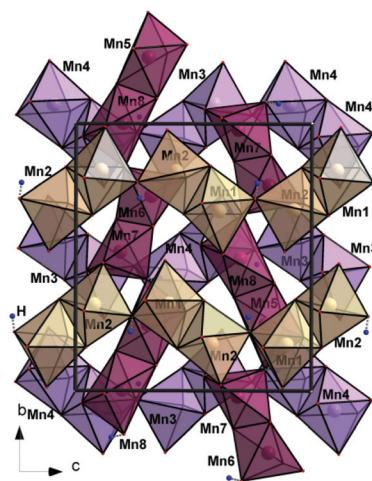


Fig. 3 Three types of chains built by Mn-centered polyhedra that form the cationic framework of the Mn<sub>2</sub>(PO<sub>4</sub>)OH crystal structure.

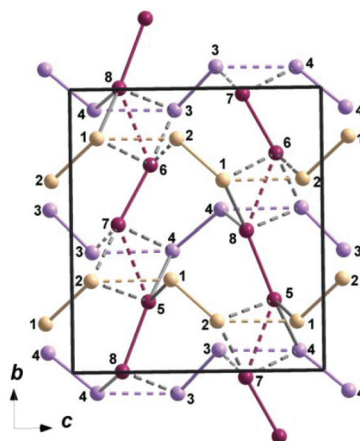


Fig. 4 Topological scheme of Mn<sup>2+</sup> cationic substructure of the synthetic triploidite. Mn<sup>2+</sup> cations are shown as spheres with numbers. Solid lines indicate edge sharing of MnO<sub>4</sub>(OH)<sub>2</sub> octahedra and MnO<sub>4</sub>(OH) bipyramids; dotted lines specify vertex sharing of neighbouring polyhedra.

interaction, while the vertex-sharing units are coupled by anti-ferromagnetic exchange interactions.<sup>24</sup> The first principles calculation of relative strength of these interactions in Mn<sub>3</sub>(PO<sub>4</sub>)<sub>2</sub> phases<sup>25</sup> in comparison to exchange interactions through [PO<sub>4</sub>]<sup>3–</sup> groups has allowed neglecting the latter.

### Physical properties

The properties of Mn<sub>2</sub>(PO<sub>4</sub>)OH, *i.e.* magnetization *M* and specific heat *C<sub>p</sub>*, were investigated using relevant options of “Quantum Design” Physical Properties Measurement System PPMS-9T in the range 2–300 K. The temperature dependence of magnetic susceptibility  $\chi = M/B$  of Mn<sub>2</sub>(PO<sub>4</sub>)OH taken at *B* = 0.1 T is shown in Fig. 5. Below room temperature, the  $\chi(T)$  dependence follows the Curie–Weiss law  $\chi = \chi_0 + C/(T - \theta)$  modified by the addition of the temperature-independent term  $\chi_0$ . The fit in the range 200–300 K with  $\chi_0 = 9.6 \times 10^{-4}$  emu mol<sup>–1</sup> yields Curie constant *C* = 8.66 emu K mol<sup>–1</sup> and

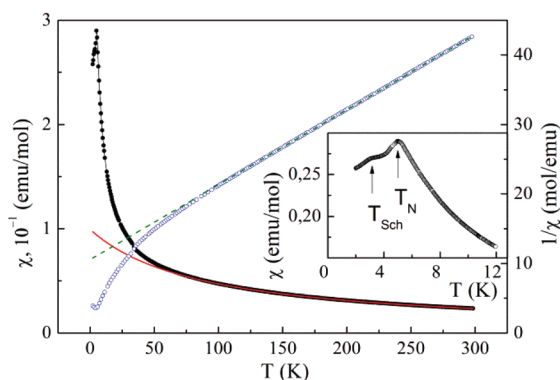


Fig. 5 The temperature dependences of both direct and inverse magnetic susceptibility in Mn<sub>2</sub>(PO<sub>4</sub>)OH taken at *B* = 0.1 T. The inset enlarges the low-temperature data.

Weiss temperature  $\Theta = -88$  K. The value of  $C$  corresponds to the squared effective magnetic moment  $\mu_{\text{eff}}^2 = 69.3\mu_{\text{B}}^2$  per formula unit, which is in good agreement with the calculated value  $\mu_{\text{eff}}^2 = ng^2S(S+1)\mu_{\text{B}}^2 = 70\mu_{\text{B}}^2$  for two magnetically active  $\text{Mn}^{2+}$  ions ( $S = 5/2$ ) per formula unit ( $n = 2$ ) and  $g$  factor  $g = 2$ . A large negative value of  $\Theta$  indicates the predominance of antiferromagnetic exchange interactions at room temperature.

At about 50 K, the  $\chi(T)$  dependence deviates upward from the extrapolation of the Curie–Weiss law which evidences the presence of ferromagnetic exchange interactions. The sharp peak at about 5 K marks the transition into a long range antiferromagnetic state. Strictly speaking, the phase transition into a magnetically ordered state is defined by singularity in  $d\chi/dT$  vs.  $T$  dependence, rather than singularity in  $\chi(T)$  dependence itself, cf. specific heat data. Below this temperature, an additional relatively broad anomaly is seen at  $T_{\text{Sch}} = 3$  K.

The field dependence of magnetization  $M(B)$  in  $\text{Mn}_2(\text{PO}_4)\text{OH}$  obtained at  $T = 2$  K is shown in Fig. 6. The magnetization increases non-monotonously reflecting tentatively the spin-flop transition marked by a sharp peak in  $dM/dB$  vs.  $B$  dependence, as shown in the inset to Fig. 6. The low value of spin-flop field  $B_{\text{SF}} = 0.9$  T is consistent with low crystalline anisotropy assumed for  $\text{Mn}^{2+}$  ions. At  $B = 9$  T, the magnetization reaches only quarter of presumed saturation magnetization of about  $10\mu_{\text{B}}$  signaling rather large values of exchange interaction parameters in  $\text{Mn}_2(\text{PO}_4)\text{OH}$ .

Further evidence of long range magnetic ordering was obtained in measurements of specific heat  $C_p$ . Its temperature dependence is shown in Fig. 7. Along with a sharp peak at  $T_N = 4.6$  K two broad Schottky-type anomalies can be seen both below and above the Néel temperature. The low temperature Schottky-type anomaly is more pronounced in  $C_p/T$  vs.  $T$  presentation of experimental data, as shown in the inset to Fig. 7. The Schottky-type anomalies in both magnetic susceptibility and specific heat could be ascribed to some non-magnetic entity (tentatively hydrogen), since it is virtually insensitive to an external magnetic field.

The magnetic entropy released at  $T < T_N$  constitutes about 20% of the total magnetic entropy  $S_{\text{magn}} = nR \ln(2S + 1)$ , where

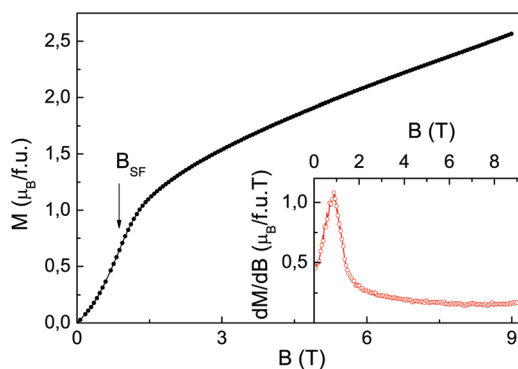


Fig. 6 The field dependence of magnetization in  $\text{Mn}_2(\text{PO}_4)\text{OH}$  at  $T = 2$  K. The arrow at  $B_{\text{SF}}$  denotes the spin-flop transition. The inset represents the  $dM/dB$  vs.  $B$  dependence.

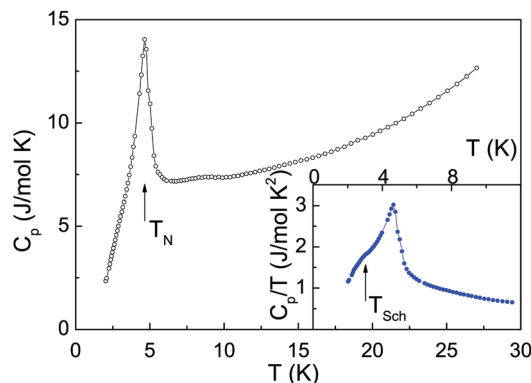


Fig. 7 The temperature dependence of specific heat in  $\text{Mn}_2(\text{PO}_4)\text{OH}$ . The inset visualizes the low temperature Schottky-type anomaly at  $T_{\text{Sch}} = 3$  K.

$R$  is the universal gas constant. This fact is in correspondence with a large value of the  $|\Theta|/T_N$  ratio indicating that magnetic entropy is released mainly at the formation of a short-range correlation regime well above the Néel temperature.

## Discussion

Despite the fact that physical properties of both natural and synthetic triplite–triploidite series of compounds are not well documented yet, it is possible to state that the magnetic  $\text{Mn}_2(\text{PO}_4)\text{OH}$  phase stands well apart from any other member of this family studied so far. Thus, both triplite  $\text{Mn}_2(\text{PO}_4)\text{F}$  and triploidite-type  $\text{Co}_2(\text{PO}_4)\text{F}$  compounds<sup>26</sup> are characterized by modest frustration strength  $|\Theta|/T_N \sim 3$ . The same refers to  $\text{Fe}_2(\text{AsO}_4)\text{F}$  with a triplite structure.<sup>27</sup>

The geometry of strongly corrugated intersecting chains of Mn-centered polyhedra in  $\text{Mn}_2(\text{PO}_4)\text{OH}$  (Fig. 3 and 4) does not presume heavy magnetic frustration, at least when the next-nearest-neighbor interactions are not taken into account. These chains are composed of edge- and vertex-sharing units, which presume an alternation of ferromagnetic and antiferromagnetic exchange interactions in accordance with Goodenough–Kanamori–Anderson rules. The competition between antiferromagnetic and ferromagnetic interactions may lead to the suppression of Weiss temperature but does not influence the Néel temperature significantly.

Therefore, the source of heavy frustration is to be seen in the arrangement of Mn-centered polyhedra along the  $a$  axis. The vertex-sharing polyhedra in the  $\text{Mn}_2(\text{PO}_4)\text{OH}$  crystal structure form the tubular fragment, as shown in Fig. 8a. Within these tubes the twisted saw tooth chain arrangement of vertex-sharing Mn-centered polyhedra can be identified, as shown in Fig. 8b. This arrangement is topologically equivalent to the ribbon cut from the highly frustrated kagomé lattice. The interactions between tubes in the  $\text{Mn}_2(\text{PO}_4)\text{OH}$  triploidite structure are provided by edge-sharing polyhedra (see Fig. 4), being tentatively ferromagnetic.



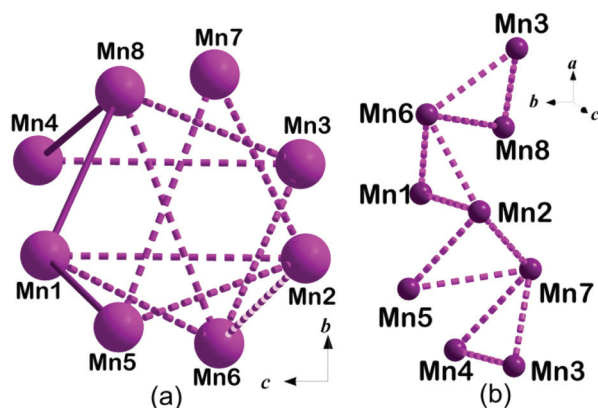


Fig. 8 Tubular structural fragment formed by  $\text{Mn}^{2+}$  cations shown along the  $a$  axis (a). The geometry of corner-sharing twisted saw tooth chain within the tube of Mn atoms (b).

Despite the fact that the  $\text{Co}_2(\text{PO}_4)\text{F}$  triploidite-type phase is isostructural to the title compound, it is not heavily frustrated. In the case of  $\text{Co}_2(\text{PO}_4)\text{F}$ , the polyhedra within the tubular fragments running along the  $a$  axis share both fluorine and oxygen vertices. The presumed difference in bonding through the overlap of  $d$  orbitals of  $\text{Mn}^{2+}$  ions with  $p$  orbitals of either  $\text{F}^-$  or  $\text{O}^{2-}$  ions lifts the frustration. The triplite,  $\text{Mn}_2(\text{PO}_4)\text{F}$ , is not frustrated, since no triangular motif can be identified in the crystal structure of this compound organized by edge-sharing octahedra.

Let us note here that phosphate mineral analogues  $\text{Co}_2(\text{PO}_4)\text{OH}$  and  $\text{Cu}_2(\text{PO}_4)\text{OH}$  can be obtained in completely different structure types, the former showing the spin-glass behavior and the latter showing properties of the spin liquid.<sup>28</sup>

## Conclusions

Summarizing, we have found that synthetic triploidite,  $\text{Mn}_2(\text{PO}_4)\text{OH}$ , is characterized by the strength of magnetic frustration comparable to that seen in geometrically frustrated spinels, pyrochlores or garnets. This three-dimensional antiferromagnetic phase possesses the high ratio of magnetic frustration due to the twisted saw tooth chain geometry of antiferromagnetically coupled polyhedra within ferromagnetically coupled tubes. The predominance of antiferromagnetic exchange interactions follows from the large negative value of Weiss temperature. At the same time, the presence of the ferromagnetic exchange interactions in the structure of  $\text{Mn}_2(\text{PO}_4)\text{OH}$  is evidenced by a strong deviation of  $\chi(T)$  from the Curie–Weiss law approaching the Néel temperature. It is definitely of interest to solve eventually the magnetic structure of this compound with quite unusual magnetic topology.

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## Notes and references

- 1 A. P. Ramires, *Annu. Rev. Mater. Sci.*, 1994, **24**, 453–480.
- 2 W. Sun, Y.-X. Huang, S. Nokhrin, Y. Pan and J.-X. Mi, *J. Mater. Chem. C*, 2016, **4**, 8772–8777.
- 3 D. Boldrin, K. Knight and A. S. Wills, *J. Mater. Chem. C*, 2016, **4**, 10315–10322.
- 4 T. Yavors’kii, M. Enjalran and M. J. P. Gingras, *Phys. Rev. Lett.*, 2006, **97**, 267203.
- 5 S. T. Bramwell and M. J. P. Gingras, *Science*, 2001, **294**, 1495–1501.
- 6 E. M. Wheeler, B. Lake, A. T. M. N. Islam, M. Reehuis, P. Steffens, T. Guidi and A. H. Hill, *Phys. Rev. B: Condens. Matter*, 2010, **82**, 140406.
- 7 C. Lacroix, P. Mendels and F. Mila, ed, *Introduction to Frustrated Magnetism*, Springer Series in Solid-State Sciences, 2011.
- 8 G. C. Lau, T. Klimczuk, F. Ronning, T. M. McQueen and R. J. Cava, *Phys. Rev. B: Condens. Matter*, 2009, **80**, 214414.
- 9 T. Armbruster, C. A. Geiger and G. A. Lager, *Am. Mineral.*, 1992, **77**, 512–521.
- 10 H. Liu, L. Yuan, S. Wang, H. Fang, Y. Zhang, C. Hou and S. Feng, *J. Mater. Chem. C*, 2016, **4**, 10529–10537.
- 11 Y. Wan and Y. B. Kim, *Phys. Rev. B: Condens. Matter*, 2016, **94**, 224401.
- 12 A. M. Hallas, J. A. M. Paddison, H. J. Silverstein, A. L. Goodwin, J. R. Stewart, A. R. Wildes, J. G. Cheng, J. S. Zhou, J. B. Goodenough, E. S. Choi, G. Ehlers, J. S. Gardner, C. R. Wiebe and H. D. Zhou, *Phys. Rev. B: Condens. Matter*, 2012, **86**, 134431.
- 13 H. Strunz and E. H. Nickel, *Strunz mineralogical tables. Chemical-structural mineral classification system*, 9th edn, Schweizerbart, Stuttgart, 2001.
- 14 Ch. Chopin, Th. Armbruster, E. Grew, A. Baronnet, C. Leyx and O. Medenbach, *Eur. J. Mineral.*, 2014, **26**, 553–565.
- 15 B. Lazic, Th. Armbruster, Ch. Chopin, E. S. Grew, A. Baronnet and L. Palatinus, *Acta Crystallogr., Sect. B: Struct. Sci.*, 2014, **70**, 243–258.
- 16 I. de Pedro, J. M. Rojo, J. Rius, O. Vallcorba, I. R. de Larramendi, J. Rodríguez Fernandez, L. Lezama and T. Rojo, *Inorg. Chem.*, 2012, **51**, 5246–5256.
- 17 L. J. Farrugia, *J. Appl. Crystallogr.*, 2012, **45**, 849–854.
- 18 *International Tables for Crystallography*, ed. E. Prince, Kluwer, Dordrecht, 3rd edn, Tables 4.2.6.8 and 6.1.14, 2004.
- 19 G. M. Sheldrick, *Acta Crystallogr., Sect. A: Fundam. Crystallogr.*, 2015, **71**, 3–8.



- 20 G. M. Sheldrick, *Acta Crystallogr., Sect. C: Cryst. Struct. Commun.*, 2015, **71**, 3–8.
- 21 I. D. Brown and D. Altermatt, *Acta Crystallogr., Sect. B: Struct. Sci.*, 1985, **41**, 244–247.
- 22 I. D. Brown, *Acta Crystallogr., Sect. A: Fundam. Crystallogr.*, 1976, **32**, 24–31.
- 23 I. D. Brown and R. D. Shannon, *Acta Crystallogr., Sect. A: Fundam. Crystallogr.*, 1973, **29**, 266–282.
- 24 J. B. Goodenough, *Magnetism and the chemical bond*, John Wiley and Sons, 1963.
- 25 O. S. Volkova, L. V. Shvanskaya, E. A. Ovchenkov, E. A. Zvereva, A. S. Volkov, D. A. Chareev, K. Molla, B. Rahaman, T. Saha-Dasgupta and A. N. Vasiliev, *Inorg. Chem.*, 2016, **55**, 10692.
- 26 M. Leblanc, I. Collin-Fèvre and G. Férey, *J. Magn. Magn. Mater.*, 1997, **167**, 71–79.
- 27 T. Berrocal, J. L. Mesa, J. L. Pizarro, M. K. Urtiaga, M. I. Arriortua and T. Rojo, *J. Solid State Chem.*, 2006, **179**, 1659–1667.
- 28 D. Karmakar and J. V. Yakhmi, *J. Phys.: Condens. Matter*, 2012, **24**, 43.

